



## Optical high-resolution spectroscopic study of $\text{Tm}^{3+}$ crystal-field levels in $\text{LiLuF}_4$

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**Abstract:** We report on the first high-resolution study of  $\text{LiLuF}_4:\text{Tm}^{3+}$ . The accurate energy level scheme of  $\text{Tm}^{3+}$  in the  $\text{LiLuF}_4$  matrix was obtained for the  ${}^3\text{H}_{6,5,4}$ ,  ${}^3\text{F}_{4,3,2}$ , and  ${}^1\text{G}_4$  multiplets. It was shown that electric-dipole transitions dominate for all the studied multiplets except the  ${}^3\text{H}_5$  one.

**Keywords:** Fourier-transform spectroscopy;  $\text{Tm}^{3+}$  crystal-field levels;  $\text{LiLuF}_4$ ; rare earths

$\text{LiYF}_4$  crystals doped with rare earth (RE) ions have proved themselves to be laser systems<sup>[1]</sup>. Optical spectroscopic study has allowed obtaining energy levels schemes of various RE ions in this matrix<sup>[1–3]</sup>.

Lithium-yttrium and lithium-lutecium double fluorides  $\text{LiRF}_4$  ( $\text{R}=\text{Y}$  or  $\text{Lu}$ ) have a scheelite ( $\text{CaWO}_4$ ) structure with the space group  $C_{4h}^6$ . Dopant  $\text{RE}^{3+}$  ions substitute for isovalent  $\text{R}^{3+}$  and occupy a single  $S_4$ -symmetry position. The crystal field (CF) levels of a  $\text{RE}^{3+}$  ion with even number of electrons (like  $\text{Tm}^{3+}$ ) are described by non-degenerate  $\Gamma_1$  and  $\Gamma_2$  irreducible representations and by doubly degenerate  $\Gamma_{34}$  one.

Recently, a novel  $\text{LiLuF}_4:\text{Tm}^{3+}$  laser has been made and investigated<sup>[4]</sup>. However, as far as we know, the data on the scheme of CF levels of the  $\text{Tm}^{3+}$  ion in this host are rare. In this work, we presented energy levels for the  $4f^{12}$  configuration of  $\text{Tm}^{3+}$  in  $\text{LiLuF}_4$  as a result of a high-resolution spectroscopic investigation of  $\text{LiLuF}_4:\text{Tm}^{3+}$  single crystals.

### 1 Experimental

In this study, we used samples of weakly doped  $\text{LiLuF}_4:\text{Tm}^{3+}$  single crystals. They were grown by the Bridgman-Stockbarger method, as described in Ref.[5]. The thulium concentration was 0.1at.%. All samples were oriented by X-ray diffraction method and cut along or perpendicularly to the crystallographic  $c$ -axis. Polarized high-resolution (up to  $0.005\text{ cm}^{-1}$ ) absorption spectra were measured in a broad spectral ( $5000\text{--}23000\text{ cm}^{-1}$ ) and temperature ( $3.5\text{--}300\text{ K}$ ) ranges using a Fourier spectrometer BRUKER IFS 125 and a Si diode or a liquid nitrogen cooled InSb detectors.

For low-temperature measurements we used a closed-cycle cryostat Cryomech ST 403.

### 2 Results and discussion

To obtain CF level scheme of the  $\text{Tm}^{3+}$  ion in  $\text{LiLuF}_4$ , we analyzed the polarized spectra of  $\text{LiLuF}_4:\text{Tm}^{3+}$  registered at two different temperatures: 4 and 70 K. The experimental data showed that the spectra in the  $\alpha$ - ( $\vec{k} \parallel \vec{c}; \vec{E}, \vec{H} \perp \vec{c}$ ) and  $\sigma$ - ( $\vec{k}, \vec{E} \perp \vec{c}; \vec{H} \parallel \vec{c}$ ) polarizations are the same for all investigated multiplets, with the exception of the  ${}^3\text{H}_5$  one. Therefore, electric-dipole transitions dominate for a majority of the multiplets, as it follows from the selection rules for the  $S_4$  site symmetry (Table 1). All three polarizations,  $\alpha$ ,  $\sigma$ , and  $\pi$  ( $\vec{k} \perp \vec{c}, \vec{E} \parallel \vec{c}, \vec{H} \perp \vec{c}$ ), differ for the  ${}^3\text{H}_5$  multiplet, manifesting that both electric-dipole and magnetic-dipole transitions are active. This fact is in good agreement with the selection rules for magnetic-dipole transitions in a free  $\text{Tm}^{3+}$  ion. Only transitions from the ground state of the  ${}^3\text{H}_6$  multiplet to the  ${}^3\text{H}_5$  one are allowed ( $\Delta J=1$ ).

There are two different ways to denote the first excited multiplet of the  $\text{Tm}^{3+}$  ion, namely,  ${}^3\text{F}_4$  or  ${}^3\text{H}_4$ . This is because of the mixing of the  ${}^3\text{F}_4$ ,  ${}^3\text{H}_4$ , and  ${}^1\text{G}_4$  states by the spin-orbit interaction. In our work, we followed the  ${}^3\text{F}_4$  notation of Jansen<sup>[2]</sup>. The absorption spectrum of  $\text{LiLuF}_4:\text{Tm}^{3+}$  in the region of the  ${}^3\text{H}_6 \rightarrow {}^3\text{F}_4$  transition is shown in Fig. 1. This multiplet splits into three  $\Gamma_1$ , two  $\Gamma_2$ , and two  $\Gamma_{34}$  levels. To identify spectral line, we used the following notations: CF levels of the ground  ${}^3\text{H}_6$  multiplet are labeled I, II, III, etc, while those in each excited multiplets A, B, C, etc, in

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the sequence of increasing energies. Transitions from the ground state are only observed in the polarized low-temperatures spectra: three lines (5596.5, 5763.1 and 5976.0  $cm^{-1}$ ) for  $\pi$ -polarization and two lines (5762.8 and 5982.9  $cm^{-1}$ ) for  $\sigma$ -polarization. Therefore, using the selection rules of Table 1 we can claim that (1) the ground state has the  $\Gamma_2$  symmetry, (2) the three levels of the  ${}^3F_4$  multiplet observed in the  $\pi$ -polarized spectrum have the  $\Gamma_1$  symmetry and the two levels observed in the  $\sigma$ -polarized spectrum have the  $\Gamma_{34}$  symmetry. The transitions from the ground state  $\Gamma_2$  to the  $\Gamma_2$  levels are allowed only as magnetic-dipole ones and are not observed in the spectra of the  ${}^3F_4$  multiplet. To obtain the positions of these levels we used the spectra registered at 70 K. Some other lines appear in them. They are shifted to 32 and 60  $cm^{-1}$  from the main lines and correspond to transitions from upper CF sublevels of the ground multiplet. First, this gives us information on two levels of the  ${}^3H_6$  multiplet which are 32 and 60  $cm^{-1}$  and have  $\Gamma_{34}$  and  $\Gamma_1$  symmetries correspondingly (see Table 1 and Fig. 1). Second, we found the positions (5837.5 and 5972.0  $cm^{-1}$ ) of the  $\Gamma_2$  levels of the  ${}^3F_4$  multiplet. A complex line shape of the highest-energy transitions in this multiplet is due to the resonant enhancement of a vibronic transition<sup>[6]</sup> and will be discussed elsewhere.

Situation is more complicated in the case of the  ${}^3H_5$  multiplet which splits into three  $\Gamma_1$ , two  $\Gamma_2$ , and three  $\Gamma_{34}$  levels. There are both electric-dipole and magnetic-dipole transitions in the spectra (see Fig. 2): five lines (8305.0, 8326.3, 8509.5, 8524.2, and 8544.0  $cm^{-1}$ ) are observed for  $\pi$ -polarization,

five lines (8288.4, 8305.0, 8509.5, 8524.2, and 8538  $cm^{-1}$ ) for  $\sigma$ -polarization and four lines (8305.0, 8509.5, 8524.2, and 8538  $cm^{-1}$ ) for  $\alpha$ -polarization. Using again the selection rules and the spectra registered at 4 and 70 K we find out that 8326.3, 8524.2, and 8544.0  $cm^{-1}$  lines correspond to the  $\Gamma_1$  levels, 8305.0, 8509.5, and 8538  $cm^{-1}$  lines are  $\Gamma_{34}$  levels. At last, very strong 8288.4  $cm^{-1}$  line observed in  $\sigma$ -polarization corresponds to the  $\Gamma_2$  level. We did not succeed in obtaining the position of the other  $\Gamma_2$  level. Here it should be said that, first, the  $S_4$  point group is a subgroup of the  $D_{2d}$  point group, and some transitions allowed in  $S_4$  are strictly forbidden in  $D_{2d}$ . If the deviation of the actual arrangement of atoms (the  $S_4$  symmetry) from the  $D_{2d}$  symmetry is small (which is the case for  $R^{3+}$  ions in  $LiRF_4$ ), some lines in the optical spectra have weak relative intensity or are absent. Second, enhanced vibronics, like as in vicinity of 5982.9  $cm^{-1}$  line (see Fig. 1), make the determination of weak transitions obscured.

In the same way we investigated the spectra of the other ( ${}^3H_4$ ,  ${}^3F_{3,2}$ , and  ${}^1G_4$ ) multiplets. Information on the CF levels of  $Tm^{3+}$  in  $LiLuF_4$  is presented in Table 2. The obtained re-

**Table 1** Selection rules for optical transitions in the  $S_4$  site symmetry (Notations  $d_i(\mu_i)$ ,  $i = x, y, \text{ or } z$  refer to the components of electric (magnetic) dipole moment)

| $S_4$         | $\Gamma_1$  | $\Gamma_2$  | $\Gamma_{34}$   |
|---------------|---|---|---|
| $\Gamma_1$    | $\mu_z(\sigma_m)$                                       | $d_z(\pi_e)$  | $d_{xy}, \mu_{xy}(\alpha_e, \sigma_e; \alpha_m, \pi_m)$ |
| $\Gamma_2$    | $d_z(\pi_e)$  | $\mu_z(\sigma_m)$                                       | $d_{xy}, \mu_{xy}(\alpha_e, \sigma_e; \alpha_m, \pi_m)$ |
| $\Gamma_{34}$ | $d_{xy}, \mu_{xy}(\alpha_e, \sigma_e; \alpha_m, \pi_m)$ | $d_{xy}, \mu_{xy}(\alpha_e, \sigma_e; \alpha_m, \pi_m)$ | $d_z, \mu_z(\pi_e, \sigma_m)$                           |

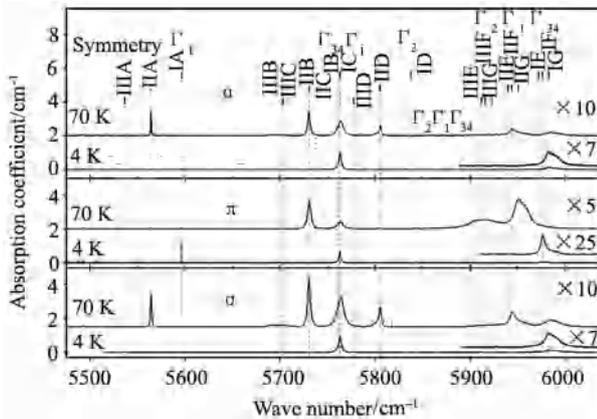


Fig. 1 Absorption spectra of  $LiLuF_4:Tm^{3+}$  in the region of the  ${}^3H_6 \rightarrow {}^3F_4$  transition

**Table 2** Experimental energy levels and their irreducible representations for  $Tm^{3+}$  in  $LiLuF_4$

| Multiplet | $\Gamma$ | Energy/ $cm^{-1}$ | Multiplet | $\Gamma$ | Energy/ $cm^{-1}$ |
|-----------|----------|-------------------|-----------|----------|-------------------|
| ${}^3H_4$ | 2        | 12598.7           | ${}^3H_6$ | 2        | 0                 |
|           | 1        | 12624.1           |           | 34       | 32.2              |
|           | 34       | 12648.3           |           | 1        | 60.0              |
|           | 1        | 12749.4           |           | 34       | 60.0              |
|           | 34       | 12842.7           |           | 1        | 60.0              |
|           | 2        | 12896.3           |           | 34       | 60.0              |
| ${}^3F_3$ | 34       | 14525.1           | ${}^3F_4$ | 34       | 5596.5            |
|           | 2        | 14547.3           |           | 1        | 5762.8            |
|           | 34       | 14602.7           |           | 34       | 5763.1            |
|           | 2        | 14605.0           |           | 1        | 5837.5            |
|           | 1        | 14612.7           |           | 2        | 5972.0            |
| ${}^3H_5$ | 34       | 15099.5           | ${}^3F_2$ | 2        | 15099.5           |
|           | 34       | 21190.3           |           | 34       | 15208.8           |
|           | 1        | 8326.3            |           | 2        | 15208.8           |
|           | 34       | 8509.5            |           | 1        | 15208.8           |
|           | 34       | 8524.2            |           | 1        | 15208.8           |
| ${}^1G_4$ | 2        | 20963.2           | ${}^3H_5$ | 2        | 8288.4            |
|           | 34       | 21190.3           |           | 34       | 8305.0            |
|           | 1        | 21280.5           |           | 1        | 8326.3            |
|           | 1        | 21302.8           |           | 34       | 8509.5            |
|           | 34       | 21564.1           |           | 1        | 8524.2            |
| ${}^3F_4$ | 2        | 21514.0           | ${}^3H_5$ | 34       | 8538.0            |
|           | 34       | 21564.1           |           | 1        | 8544.0            |
|           | 1        | 21564.1           |           | 34       | 8538.0            |
|           | 1        | 21564.1           |           | 1        | 8544.0            |
|           | 1        | 21564.1           |           | 34       | 8538.0            |

