

Upconversion of Er^{3+} Ions in $\text{LiKGdF}_6:\text{Er}^{3+}$, Dy^{3+} Single Crystal Produced by Infrared and Green Laser

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Received 10 June 2006; revised 17 October 2006

Abstract: The upconversion fluorescence of Er^{3+} ions in $\text{LiKGdF}_6:\text{Er}^{3+}$, Dy^{3+} single crystal was studied under 785, 514.5, and 980 nm laser excitations. With the laser excitations of at 785 nm, strong green (centered at 543 nm) upconversion emissions, as well as weak red (651 nm), violet (406 nm), and blue (470 nm) upconversion emissions were obtained. With 514.5 nm laser excitations, violet (406 nm) and blue (470 nm) upconversion emissions were observed. Under the 980 nm laser excitation, strong green (543 nm) and weak red (651 nm) emissions were also observed. The laser power dependency of the upconversion emissions was investigated to understand the upconversion mechanism. The critical cross absorption (CCA) and the energy transfer (ET) processes were discussed as the possible mechanisms for all upconversion emissions.

Key words: upconversion; $\text{LiKGdF}_6:\text{Er}^{3+}$; excited state absorption; energy transfer; rare earths

CLC number: O482.31; O614.33 Document code: A Article ID: 1002-0721(2006)06-0740-05

In the past decades, the upconversion of infrared light to visible light of different materials, such as single crystals, optical fibers, and glasses, has been investigated extensively for a wide range of applications, including all-solid compact laser devices operating in the violet-blue-green region, infrared quantum number detectors, transceivers systems, and fluorescent labels for sensitive detection of biomolecules and optical data storage^[1–3]. More recently, exploration on infrared or visible to violet and ultraviolet (UV) upconversion fluorescence has become considerably more significant because it may be useful for new visible all-solid-state lasers operating in the violet and UV region^[4]. In the upconversion process, two (or more) low-energy photons from the excitation source are converted into one

photon with higher energy.

Since the intermediate levels ${}^4\text{I}_{13/2}$ (around 12491 cm^{-1}) and ${}^4\text{I}_{15/2}$ (around 10275 cm^{-1}) of Er^{3+} ions can be conveniently populated by commercial low-power high-power near-infrared laser diodes, Er^{3+} ion doped materials are well suited for upconversion^[5,7–12]. Most upconversion studies of Er-doped materials were focused on infrared laser excitation.

However, investigation on violet and blue upconversion under visible laser excitation is very scarce^[4,7,8,13]. This upconversion phenomenon is yet to be explained. Fortunately, the ${}^4\text{F}_{9/2}$ level, which has a longer lifetime than that of the other states in the visible range, provides an available intermediate state for such exploration^[7,8].

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Foundation Item: Project supported by a grant from Department of Education of Zhejiang Province (2006096)

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It is well established that the luminescence efficiency of rare earth ions is governed principally by the nonradiative process of materials, and the highest phonon frequencies of the host materials are responsible for the nonradiative relaxation rate^{20,21}. Therefore, by selecting a host matrix that has lower phonon energy, high upconversion efficiency can be obtained. Fluoride crystals have been demonstrated as one of the best host materials, benefiting from their low phonon energy, wide transparent region, and their good physical and chemical stability^{22, 26, 28}.

In this article, the upconversion fluorescence of Er^{3+} ions in $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal was studied under 785, 514.5 and 980 nm laser excitation. The laser power dependence of the upconverted emission was investigated to understand the upconversion mechanism. The excited state absorption (ESA) and the energy transfer (ET) processes were discussed as the possible mechanisms for all upconversion emissions.

1 Experimental

LiKGdF_6 single crystal co-doped with Er^{3+} (2%) and Dy^{3+} (0.4%) was prepared using the hydrothermal method. The synthesis process was similar to that given in Ref. [24].

The upconversion spectra, excited by 785 and 514.5 nm laser, were measured and analyzed using a Jobin-Yvon LABRAM-HR laser micro-Raman spectroscopy system. For temperature dependent studies, samples were placed in a circulating liquid nitrogen cryostat and the temperature was set between 90 and 300 K.

Green and red upconversion emissions, excited by a 980 nm diode laser, were measured by a Jobin-Yvon HRD1 double monochromator and were detected using a Hamamatsu R456 photomultiplier. The signal was analyzed by the EG&G 7265 DSF lock-in amplifier and was stored in the computer memory.

2 Results and Discussion

2.1 Pumping with 785 nm laser

Fig. 1(a) shows the room temperature upconversion spectra of $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal under the 785 nm laser excitation. The upconversion spectra show bands centered at 522 nm (green), 543 nm (green), and 651 nm (red), which can be attributed to the $^3\text{H}_{11/2}-^4\text{I}_{15/2}$, $^3\text{S}_2-^4\text{I}_{15/2}$, and $^4\text{F}_{9/2}-^4\text{I}_{15/2}$ transitions of Er^{3+} ions, respectively. In addition to these bands, very weak emissions at 406 nm (violet) and 570 nm (blue) are also observed and attributed to the $^3\text{P}_2-^4\text{I}_{15/2}$ and $^3\text{P}_0-^4\text{I}_{15/2}$ transitions of Er^{3+} ions, respectively²⁶.

To understand the mechanisms of the observed upconversion luminescence, the upconverted luminescence intensity I of these transitions was measured as a function of the pump power P . In the upconversion process, I is proportional to the n th power of P , i.e., $I \propto P^n$ (1)

where, n is the number of pump photons required to populate the emitting state. A plot of $\lg I$ versus $\lg P$ yields a straight line with slope n . The power dependence of I is shown in Fig. 2(a, b). The slopes n obtained were 1.93, 1.97, and 1.69 for 522, 543, and

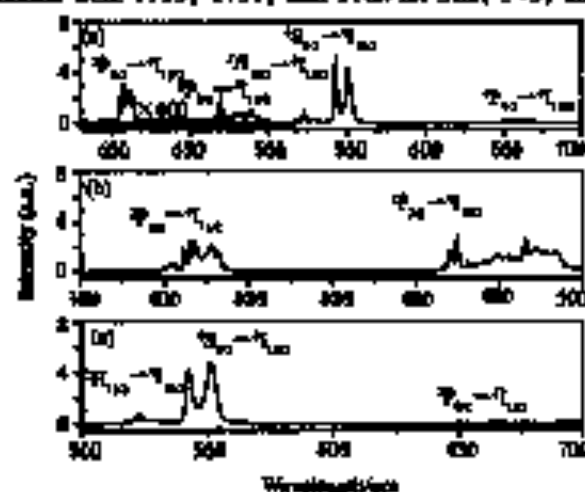


Fig. 1 Room temperature upconversion spectra of $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal
(a) $\lambda_{\text{pump}} = 785 \text{ nm}$; (b) $\lambda_{\text{pump}} = 514.5 \text{ nm}$; (c) $\lambda_{\text{pump}} = 980 \text{ nm}$

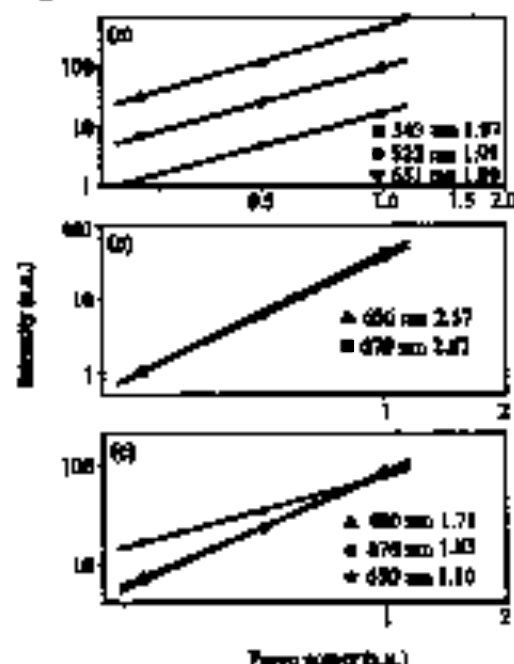


Fig. 2 Dependence of the upconversion emission intensity on excitation power for $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal
(a) $\lambda_{\text{pump}} = 785 \text{ nm}$ (for green and red emissions); (b) $\lambda_{\text{pump}} = 785 \text{ nm}$ (for violet and blue emission); (c) $\lambda_{\text{pump}} = 514.5 \text{ nm}$

651 nm emissions, respectively. The results show that two-photon processes contribute to the green and red emissions. For 405 and 470 nm emissions, the n values obtained were 2.57 and 2.57, respectively. This result strongly indicates that three-photon processes contribute to the violet and blue emissions.

The mechanisms of visible emissions of Er^{3+} ions under excitation into the ${}^4\text{I}_{13/2}$ level have been well discussed in the literature, which involve two well-known upconversion mechanisms: (1) the excited state absorption (ESA), and (2) the energy transfer (ET). The upconversion paths in $\text{LiKGdF}_4:\text{Er}^{3+}$, Dy^{3+} are now analyzed. Fig. 3(a) provides the energy level diagram of Er^{3+} ions and the possible upconversion mechanisms under 785 nm laser excitation^{26,100}.

When the ${}^4\text{I}_{13/2}$ level of Er^{3+} is directly excited with 785 nm laser, the initial ${}^4\text{I}_{13/2}$ level population increases to long lifetimes ${}^3\text{H}_{11/2}$ and ${}^3\text{I}_{13/2}$ levels. The ions in ${}^3\text{I}_{13/2}$ level sequentially absorb another 785 nm photon and are raised to ${}^2\text{F}_{5/2,3/2}$ levels (labeled as ESA1). The ions in ${}^4\text{I}_{13/2}$ level can also absorb another 785 nm photon and are raised to the ${}^2\text{H}_{11/2}$ level (ESA2). The ions in ${}^2\text{F}_{5/2,3/2}$ and ${}^2\text{H}_{11/2}$ levels undergo multi-photon relaxation to ${}^3\text{H}_{11/2}$ and ${}^3\text{S}_2$ levels and provide green emissions. There is a possible ET route that can also populate these luminescent levels. An excited ion relaxes from the ${}^3\text{I}_{13/2}$ level to the ${}^3\text{I}_{15/2}$ level nonradiatively and transfers the emission energy to the neighboring ion in the same level, thus promoting the latter to the ${}^2\text{F}_{7/2}$ level: ${}^3\text{I}_{13/2} + {}^3\text{I}_{15/2} \rightarrow {}^3\text{I}_{15/2} + {}^2\text{F}_{7/2}$ (ET1). It must be mentioned that the ET1 process is a phonon-assisted process (absorbing phonon) with an energy mismatch of about 200 cm^{-1} . The ${}^2\text{F}_{7/2}$ level can be possibly pumped via a nonradiative relaxation through the ${}^3\text{S}_2$ excited state at moderate rate.



Fig. 3 Energy level diagram of Er^{3+} ions in $\text{LiKGdF}_4:\text{Er}^{3+}$, Dy^{3+} single crystal and upconversion mechanisms (a) $\lambda_{\text{ex}} = 785 \text{ nm}$, (b) $\lambda_{\text{ex}} = 514.5 \text{ nm}$, (c) $\lambda_{\text{ex}} = 980 \text{ nm}$

As stated above, the power dependence of the 405 (${}^2\text{P}_{3/2} \rightarrow {}^4\text{I}_{13/2}$ transition) and 470 nm (${}^2\text{P}_{1/2} \rightarrow {}^4\text{I}_{13/2}$ transition) emissions indicates that a three-photon process is involved. The ${}^2\text{P}_{3/2}$ state is populated by the following mechanism: When the ${}^3\text{S}_2$ level is populated, the ions in the ${}^3\text{S}_2$ level can absorb another 785 nm photon and are raised to the ${}^2\text{P}_{3/2}$ level (ESA3).

To obtain a first idea on the nature of the upconversion mechanisms, the effect of temperature on the upconversion luminescence has been investigated as shown in Fig. 4. Two phenomena were observed in this figure.

One is that almost no spectral line around 322 nm (${}^2\text{H}_{11/2} \rightarrow {}^4\text{I}_{13/2}$ emission) could be observed in Fig. 4 (3). This is because the energy gap between the ${}^2\text{H}_{11/2}$ and ${}^3\text{S}_2$ levels is about 770 cm^{-1} and the population of Er^{3+} ions in these two levels depends on the thermal equilibrium distribution. At room temperature, these two levels have almost the same population, whereas at 90 K, the population of the ${}^2\text{H}_{11/2}$ level is very low.

The other phenomenon is that the intensities of green lines corresponding to the ${}^3\text{S}_2 \rightarrow {}^4\text{I}_{13/2}$ transition and the intensities of red lines decrease with the decreasing temperature. Both a phenomenon is caused by the fact that the ET1 process is a phonon-assisted process (absorbing phonon) and the energy mismatch is about 200 cm^{-1} . At low temperature (90 K), the rate of the ET1 process decreases considerably and hence the emission intensity decreases.

2.2 Pumping with 514.5 nm laser

Fig. 1(b) shows the room temperature upconversion spectra of $\text{LiKGdF}_4:\text{Er}^{3+}$, Dy^{3+} single crystal under 514.5 nm laser excitation. The upconversion spectra show bands centered at 405 and 470 nm, which are attributed to the ${}^2\text{P}_{3/2} \rightarrow {}^4\text{I}_{13/2}$ and ${}^2\text{P}_{1/2} \rightarrow {}^4\text{I}_{13/2}$

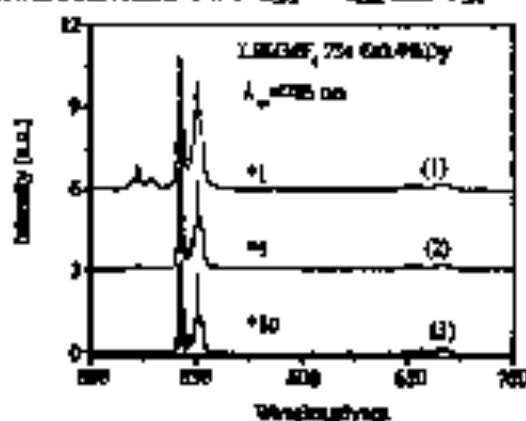


Fig. 4 Upconversion luminescence spectra of $\text{LiKGdF}_4:\text{Er}^{3+}$, Dy^{3+} single crystal at different temperatures (1) 300 K, (2) 120 K and (3) 90 K

transitions of Er^{3+} ions, respectively¹⁷. The bands in the region of 477–500 nm may arise from the ${}^4\text{F}_{7/2} \rightarrow {}^4\text{I}_{13/2}$ transition of the Er^{3+} ion.

Fig. 2(c) shows the intensity dependence of the upconversion emissions versus the pump power. The slope n for 486, 477, and 496 nm was 1.71, 1.83, and 1.10, respectively. This result indicates that two-photon processes contribute to the emission from the ${}^2\text{F}_{5/2}$ level and then the one-photon process pumps the ions at the ${}^4\text{F}_{7/2}$ level. The energy mismatch between the emission at 486 nm and the excitation laser wavelength 514.5 nm amounts to 1140 cm^{-1} . Taking into account the thermal population (at room temperature) of Stark components of the ${}^4\text{I}_{13/2}$ level, Er^{3+} ions can be excited from the ground state to the ${}^4\text{F}_{7/2}$ level by absorbing one 514.5 nm photon with the assistance of phonons.

Fig. 3(b) depicts the probable upconversion mechanism¹⁷. Under 514.5 nm laser excitation, Er^{3+} ions are first excited from ground state to ${}^2\text{H}_{11/2}$ or ${}^4\text{F}_{7/2}$ levels by absorbing one pump photon, and then relax quickly to the lower long lifetime level ${}^2\text{S}_{3/2}$. The ions in the ${}^2\text{S}_{3/2}$ level sequentially absorb another 514.5 nm photon and are raised to the ${}^4\text{F}_{7/2}$ level (ESA4). There is a possible ET route that can also populate these luminescent levels. An excited ion relaxes from the ${}^2\text{S}_{3/2}$ level to the ${}^4\text{I}_{13/2}$ level nonradiatively and transfers the excitation energy to the neighboring ion in the same level, promoting it to the ${}^4\text{G}_{7/2}$ levels: ${}^2\text{S}_{3/2} + {}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{13/2} + {}^4\text{G}_{7/2}$ (ET2). Finally, the ions relax to the ${}^2\text{F}_{5/2}$ level through multi-phonon nonradiative processes, and then provide the upconversion luminescence.

The low temperature upconversion spectra of $\text{LiK-GdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal excited by 514.5 nm laser (not shown here) indicates that the intensities of emissions from the ${}^2\text{F}_{5/2}$ level (486 and 470 nm) increase with the decreasing temperature and almost no excitation from the ${}^2\text{F}_{5/2}$ level was observed at 90 K. The former is caused by the fact that at low temperature, there is room population to the ${}^2\text{F}_{5/2}$ level because of the low nonradiative rate. The latter can be explained that the ${}^4\text{F}_{7/2}$ level is populated via absorbing one 514.5 nm photon from the Stark components of the ground state.

3.3 Pumping with 980 nm laser

Fig. 1(c) shows the room temperature upconversion spectra of $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal excited by 980 nm laser. The bands centered at 522, 551, and 651 nm are assigned to the ${}^4\text{H}_{11/2} \rightarrow {}^4\text{I}_{13/2}$, ${}^2\text{S}_{3/2} \rightarrow {}^4\text{I}_{13/2}$ and ${}^4\text{F}_{7/2} \rightarrow {}^4\text{I}_{13/2}$ transitions of Er^{3+} ions, re-

spectively, which is identical to Fig. 1(a).

The intensity dependence of the upconversion emissions versus the pump power (figure not shown here) shows the slope n for 522, 551, and 651 nm is 2, which indicates that two-photon processes contribute to the green and red emissions.

The probable upconversion mechanisms are presented in Fig. 3(c)^{17,18}. Under 980 nm excitation, Er^{3+} ions are excited from the ground state to the ${}^4\text{I}_{13/2}$ level by absorbing one 980 nm laser photon (ESA5). Ions in the ${}^4\text{I}_{13/2}$ level sequentially absorb another 980 nm photon and are raised to the ${}^4\text{F}_{7/2}$ level. Ions in the ${}^4\text{F}_{7/2}$ level undergo multi-phonon relaxation to luminescent levels ${}^2\text{H}_{11/2}$ and ${}^2\text{S}_{3/2}$ and provide green emission from these two levels. In addition, the contribution of the ET process cannot be rejected. The luminescent level ${}^2\text{F}_{5/2}$ can be populated via nonradiative relaxation from the ${}^2\text{S}_{3/2}$ excited state.

3 Conclusion

Violet, blue, green, and red upconversion luminescence emissions of Er^{3+} in $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal were observed under 783, 514.5, and 980 nm laser excitations. The excited state absorption and the energy transfer process were discussed as possible mechanisms for all the upconversion emissions. In conclusion, the $\text{LiKGdF}_6:\text{Er}^{3+}, \text{Dy}^{3+}$ single crystal is a very promising crystal material to obtain violet or visible upconversion emission.

Acknowledgments: The authors are grateful to Dr. V. G. Misher and Dr. N. M. Koshchev for providing the sample.

References:

- [1] Kojima F, Motomura F, Drapeau T, et al. Green upconversion emission uses $\text{Er}^{3+}:\text{LiYF}_6$ laser at room temperature [J]. *Applied Physics Letters*, 1994, 65: 1883.
- [2] Valentin R, Wang Q S, Qiu H U. Upconversion luminescence in Yb^{3+} doped BaMgCl_2 : Spectroscopy, detection, and applications [J]. *Journal of Chemical Physics*, 2002, 116: 3190.
- [3] Qiu C, Huang S, Peng Y, et al. Power scaling of Yb^{3+} doped ZBLAN fiber structures fiber lasers: Modeling and experiment [J]. *Applied Physics B: Laser and Optics*, 2005, 82: 65.
- [4] Zhang J, Dai S, Wang G, et al. Investigation on upconversion luminescence in $\text{Er}^{3+}/\text{Yb}^{3+}$ codoped alkoxide glasses and fibers [J]. *Physics Letters, Section A: General, Atomic and Solid State Physics*, 2006, 343: 409.
- [5] Liu H, Marshall G, Wang S, et al. Optical transitions and visible upconversion in Er^{3+} doped alkoxide tellurite glass [J]. *Journal of Applied Physics*, 2003, 95: 140.
- [6] Cao Y, Wang J, Chen Y. UV upconversion luminescence of Er^{3+} and Er^{3+} ions in fluorozirconate glasses

- [7] Gao H, Li Y, Wang L, et al. Blue upconversion of cubic $\text{Ca}_2\text{O}_3:\text{Er}^{3+}$ prepared by green laser [J]. *Journal of Alloys and Compounds*, 2009, 479: 23.
- [8] Guo U, Dong H, Xia M, et al. Visible upconversion in van with lan-doped Ca_2O_3 nanocrystals [J]. *Journal of Physical Chemistry B*, 2004, 108: 19239.
- [9] Guo H, Dong H, Xia M, et al. Green and red upconversion luminescence in Er^{3+} -doped and $\text{Er}^{3+}/\text{Tb}^{3+}$ -co-doped BaTiO_3 ultrathin powders [J]. *Journal of Alloys and Compounds*, 2006, 415: 200.
- [10] Guo H, Zhang W, Xia M, et al. Structure property and visible upconversion of Er^{3+} doped Ca_2O_3 nanocrystals [J]. *Journal of Rare Earths*, 2004, 22: 261.
- [11] He H, Shi S, Xu S, et al. Upconversion emission in Er^{3+} doped novel lamellar glass [J]. *Journal of Rare Earths*, 2005, 23: 331.
- [12] Yang J, Zhang L, Wen L, et al. Optical modified and upconversion luminescence of $\text{Er}^{3+}/\text{Tb}^{3+}$ -co-doped tellurite modified silica glasses [J]. *Journal of Applied Physics*, 2004, 95: 3620.
- [13] Arvid F. Upconversion and anti-stokes process with F and G lines in solids [J]. *Chemical Reviews*, 2004, 84: 190.
- [14] Wang D, Yin M, Shi B, et al. Upconversion fluorescence of Er^{3+} ions: impurity free and Raman study in $\text{K}_2\text{TiF}_6:0.1\text{ mol}\% \text{Th}^{4+}$ single crystal [J]. *Journal of Alloys and Compounds*, 2004, 388: 237.
- [15] Fan Z, Wang J, Qian Z, et al. Preparation process and upconversion luminescence of Er^{3+} -doped glass ceramic containing $\text{Ba}_2\text{La}_2\text{F}_7$ nanocrystals [J]. *Journal of Physical Chemistry B*, 2006, 110: 3890.
- [16] Han K, Kampa K, Ladd M U, et al. Highly efficient multicolor upconversion emission in nanosized colloids of lanthanide-doped BaF_2 nanocrystals [J]. *Advanced Materials*, 2004, 16: 2102.
- [17] Peter A, Fedor C B, Kasper H, et al. Fluorescence upconversion properties of Er^{3+} -doped TiO_2 and BaTiO_3 nanocrystallites [J]. *Chemistry of Materials*, 2009, 15: 3690.
- [18] Sun H Y, Wen L, Fan Z C, et al. Infrared frequency upconversion fluorescence emission of $\text{Er}^{3+}/\text{Tb}^{3+}$ -co-doped oxytellurite germanate glass [J]. *Journal of Alloys and Compounds*, 2006, 414: 162.
- [19] Vozzov F, Boyer J C, Capobianco J A, et al. Effect of Th^{4+} doping on the upconversion emission in nanocrystalline $\text{Y}_2\text{O}_3:\text{Er}^{3+}$ [J]. *Journal of Physical Chemistry B*, 2005, 109: 1107.
- [20] Zhou A, Hou Y, Pei Z, et al. Upconversion luminescence of $\text{YLF}_6:\text{Er}^{3+}$ synthesized by hydrothermal method [J]. *Journal of Rare Earths*, 2005, 23: 857.
- [21] Tsvetkov P A. Luminescence and collective decay of Er^{3+} in $\text{Ca}_2\text{NaErCl}_6$ [J]. *Soviet Physics*, 1988, 63: 265.
- [22] Xu H, Jiang Z. Ultraviolet and violet upconversion luminescence in Er^{3+} -doped yttrium aluminum garnet crystals [J]. *Physical Review B-Condensed Matter and Materials Physics*, 2003, 68: 351091.
- [23] Peter A, Fedor C B, Kasper H, et al. Effect of crystal sites on upconversion luminescence in $\text{Er}^{3+}:\text{BaTiO}_3$ nanocrystals [J]. *Applied Physics Letters*, 2007, 91: 286.
- [24] Yin M, Li Y, Dong H, et al. Spectroscopic studies and crystal field calculation for Er^{3+} in single crystal Ca_2VF_6 [J]. *Journal of Alloys and Compounds*, 2005, 393: 95.